

# Lethal Mutants and Truncated Selection Together Solve a Paradox of the Origin of Life

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#### **Abstract**

**Background:** Many attempts have been made to describe the origin of life, one of which is Eigen's cycle of autocatalytic reactions [Eigen M (1971) Naturwissenschaften 58, 465–523], in which primordial life molecules are replicated with limited accuracy through autocatalytic reactions. For successful evolution, the information carrier (either RNA or DNA or their precursor) must be transmitted to the next generation with a minimal number of misprints. In Eigen's theory, the maximum chain length that could be maintained is restricted to 100-1000 nucleotides, while for the most primitive genome the length is around 7000-20000. This is the famous error catastrophe paradox. How to solve this puzzle is an interesting and important problem in the theory of the origin of life.

Methodology/Principal Findings: We use methods of statistical physics to solve this paradox by carefully analyzing the implications of neutral and lethal mutants, and truncated selection (i.e., when fitness is zero after a certain Hamming distance from the master sequence) for the critical chain length. While neutral mutants play an important role in evolution, they do not provide a solution to the paradox. We have found that lethal mutants and truncated selection together can solve the error catastrophe paradox. There is a principal difference between prebiotic molecule self-replication and protocell self-replication stages in the origin of life.

**Conclusions/Significance:** We have applied methods of statistical physics to make an important breakthrough in the molecular theory of the origin of life. Our results will inspire further studies on the molecular theory of the origin of life and biological evolution.

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#### Introduction

The puzzle about the origin of life has attracted the attention of curious minds from the dawn of human civilization. Since the development of molecular biology, it has been known that the information carriers of living organisms, from humans to bacteria and viruses, are DNA and RNA. An essential step in solving the puzzle about the origin of life at a molecular level is to understand the replication and evolution of information carriers. For this purpose, Eigen [1,2] proposed a cycle of autocatalytic reactions.

Primordial life molecules are replicated through autocatalytic reactions with a limited accuracy, i.e. an error rate in the order of  $0.001 \sim 0.01$  [3,4]. For successful evolution, genetic information must be transmitted to the next generation with a minimal number of misprints. With an error rate in the order of  $0.001 \sim 0.01$  [3,4], the maximum length of information carrier that could be maintained is estimated in Eigen's theory [1,2] to be 100-1000. For the most primitive genome, the length is estimated by Gil *et al* [5] to be around 20000 nucleotides, and by Kun *et al* [6] to be around 7000-8000 nucleotides. The former includes the core bacterial gene set, and the latter includes only the key information

carrier. The big gap between 100-1000 and 7000-8000 is the famous error catastrophe paradox.

In this paper, we use methods of statistical physics to solve the error catastrophe paradox by carefully analyzing the roles of neutral networks, in which each mutant in the network has about the same reproduction rate as the master sequence [6,7], lethal mutants, i.e., mutants with a zero reproduction rate [8,9], and truncated selection, i.e., when the mutants with Hamming distances from the master sequence larger than a critical value n have a zero reproduction rate [10–12]. We calculate the impact of neutrality for the neutral thick hierarchic tree, and derive simple exact formulae for the case of the neutral network-like fitness landscape. The importance of neutral network-like fitness landscapes is widely known [13–15].

The neutrality phenomenon with perfect (the neutral mutants have exactly the same fitness as the master sequence) and extensive neutrality has been considered in [6,16,17], and a large increase in the mean fitness due to neutrality has been found. However, mutants with both perfect and extensive neutrality are not realistic and cannot be found in real biological systems. A more realistic case is imperfect extensive neutrality, to be discussed below.

In the present paper, we consider different versions of neutrality: the neutral network-like fitness landscape, which is very popular among biologists, and the more involved thick hierarchic tree landscape, to be defined later. In both cases, we derive analytical results for the mean fitness and for the probabilities of the main sequences appearing. Our result for the mean fitness of the neutral network is consistent with the rigorous result of Nimwegen *et al* [14]. In all cases we consider, the modification of the mean fitness due to neutrality involves a small factor of the order  $1/\sqrt{L}$ . The corresponding change in the critical chain length is negligible in solving the error catastrophe paradox. We will discuss the extensive neutrality [6,16,17] in the subsection **Extensive neutrality**, below. According to our analysis of the experimental data, the increase in mean fitness is also negligible for the observed case of imperfect extensive neutrality.

Applying statistical physical methods used in earlier papers [18–28] to the Eigen model with lethal mutants and truncated selection, we analyze the paradox of the origin of life. We find that the combined action of lethal mutants and truncated selection makes the error threshold reach the required genome length for the origin of life and thus solves the paradox of the origin of life.

Here we use the concept from the statistical physics of spin models [29] to review briefly Eigen's theory of the cycle of autocatalytic reactions [1,2].

The genetic information of a biological system is stored in the DNA or RNA sequence. Eigen used models similar to the one-dimensional Ising model [29] with L spins to represent DNA or RNA of L bases, and considered the time evolution of the probability distribution  $p_i$ ,  $0 \le i \le M-1$ , of  $M=2^L$  spin configurations  $S_i \equiv (s_1^{(i)}, \ldots, s_L^{(i)})$  corresponding to M DNA or RNA sequences, with +1 "spin" representing purines (R) and -1 "spin" representing pyrimidines (Y) in a sequence. Every sequence  $S_i$  is assigned a value of the fitness function,  $r_i$ . The number  $r_i$  represents the reproduction rate of  $S_i$ .

In the simplest case of the single-peak fitness function, there is only one peak configuration or master sequence, say  $S_0$ , which has the largest value of fitness function so that  $r_0 = A > 1$ , and  $r_i = 1$  for  $i \neq 0$ , as shown in Fig. 1. Configuration  $S_0$  can be chosen to be  $(1,1,\ldots,1)$ , i.e., all spins take +1, without the loss of generality. The j-th sequence  $S_j$  can change into the i-th sequence  $S_i$  via mutation. The Hamming distance between configurations  $S_i$  and  $S_k$ , i.e., the number of minimal mutation flips from  $S_k$  to  $S_i$ , is denoted by d(k,i). In the truncated selection, the fitness function is zero after some Hamming distance from the master sequence. A typical example is shown in Fig. 2.

In Eigen's theory [1,2], an information carrier reproduces at a certain rate  $r_i$ , producing offspring of the parental type with the

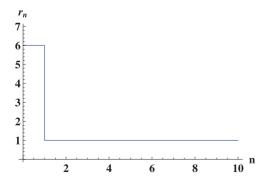


Figure 1. Fitness  $r_n$  versus Hamming distance n from the peak sequence for the single-peak fitness landscape. doi:10.1371/journal.pone.0021904.q001

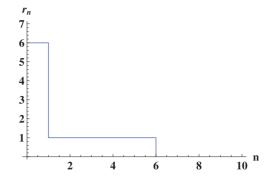


Figure 2. Fitness  $r_n$  versus Hamming distance n from the peak sequence for the truncated fitness landscape. doi:10.1371/journal.pone.0021904.g002

probability  $Q_{ii}$  and offspring of the mutant type  $S_k$   $(k \neq i)$  with the probability  $Q_{ki}$ . The probabilities  $p_i$  for different types (sequences)  $S_i$ , i = 0, ..., M-1, satisfy the set of equations

$$\frac{dp_i}{dt} = \{Q_{ii}r_i - \sum_{k} r_k p_k(t)\} p_i(t) + \sum_{k \neq i} Q_{ik} r_k p_k(t).$$
 (1)

Here,  $p_i$  satisfy the normalization condition  $\sum_{i=0}^{M-1} p_i = 1$ ; the elements of the mutation matrix are  $Q_{ki} = q^{L-d(k,i)}(1-q)^{d(k,i)}$ ; q is the probability of errorless replication per nucleotide. The diagonal terms of the mutation matrix are  $Q_{ii} = q^L \equiv Q \equiv e^{-\gamma}$ , where  $\gamma = -L \ln(q) \approx L(1-q)$  is the parameter of mutation in the Eigen model. Two sequences  $S_i$  and  $S_k$  are neighbors if and only if d(k,i)=1. For i=0, the second term on the right-hand side of Eq. (1) represents the back mutation from mutants to the master sequence.

For the single peak fitness landscape with  $r_0 = A$  and  $r_i = 1$ , for  $1 \le i \le M - 1$  (Fig. 1), Eigen derived the following restriction for the length of genome [1,2]

$$1 < QA \quad \text{or} \quad L < \frac{\ln(A)}{1 - q} = L_{max}, \tag{2}$$

for (1-q) < <1, where  $L_{max}$  is the maximal allowed genome length. If we neglect with a 1/L accuracy a small contribution from the second term in the right-hand side of Eq. (1), i.e., from back mutations, we can easily show that the steady state probability  $p_0$  for the master sequence and the mean fitness R are given by

$$p_0 = \hat{p}_0(A, Q) \equiv \frac{QA - 1}{A - 1}, \quad R \equiv \sum_{i=1}^{M} p_i r_i \approx QA.$$
 (3)

We can also use  $p_0$  and R of Eq. (3) to derive Eq. (2) from the condition  $p_0 > 0$  or R > 1. In Ref. [22], we have derived Eq. (2) as a special case of the Eigen model with a general fitness function and degradation rates.

The error rate (1-q) has a value between 0.01 and 0.001 [3,4], restricting the length of early information carriers to some number between 100 and 1000 for  $\ln A \approx 1$ , which is much smaller than 7000 or 8000 for the most primitive genome, as estimated by Kun *et al* [6], mentioned above. This is the famous error catastrophe paradox.

One hope of increasing the information content has been connected with the idea of neutral network-like fitness landscapes [6,17]. However, the quantitative impact of this phenomenon has not been rigorously investigated. Truncated fitness landscapes have been discussed in [10] with regard to Muller's ratchet.

Summers and Litwin [11] investigated the extreme form of truncated fitness with only one Hamming class for viable mutants, and claimed the absence of an error-threshold relation in virus populations. In this paper, we use methods developed in [22] to solve analytically the model with lethal mutations and truncated fitness, and find that lethal mutants and truncated selection together can solve the paradox of the origin of life.

#### Results

# Neutral Landscapes and Critical Chain Length

**Neutral network.** If two neighboring sequences (also called "nodes") have almost the same reproduction rate, the mutation from one node to another barely changes the reproduction rate. Such a mutation is called neutral mutation. In a neutral network, every node of the network has almost the same reproduction rate as the master sequence and every node in the network can be connected to the master sequence via a series of neutral mutations. In the neutral network, the number of neighboring nodes of a given node  $S_i$  is called the connectivity or degree of that node and is denoted by  $g_i$ . The probability of  $S_i$  appearing is  $p_i$ . The total probability P of the neutral sequences is  $P = \sum_{i} p_{i}$ , where the summation is over all nodes in the neutral network. The mean degree is given by  $\bar{g} = \sum_i g_i p_i / P$ . Now we have the fitness A > 1on the neutral network and fitness 1 outside the network.

In the infinite genome length limit  $(L \rightarrow \infty)$  the principal term for the mean fitness of the neutral networks and the total probability P of the neutral sequences are expected to be very close to those for the single peak fitness given by Eq. (3), hence

$$P - \hat{p}_0(A, Q) \ll 1, \quad k \equiv R - QA, \quad k \ll 1.$$
 (4)

Here,  $R \equiv \sum_i r_i p_i$  is the mean fitness computed over all sequences. The error threshold is defined from the condition QA + k = R = 1.

We consider the case that in the network there is a node with a maximal degree (also called the "0-th sequence"), for which the degree  $g \equiv g_0 > 1$ , and other nodes have small degree,  $\sim 1$ . Having a large parameter g, we can solve the evolution problem in the neutral network within  $1/\sqrt{g}$  accuracy. We assume, and our calculations confirm this conjecture, that the impact of neutrality on the mean fitness must be defined by the largest degree of the neutral network. We denote by  $p_0$  the probability of having a sequence with g neutral neighbors, and by  $p_1$  the probability of any of these g neutral sequences. Consider the steady state solution of Eq. (1). With the accuracy 1/g, we obtain the system of equations for  $p_0$  and  $p_1$ :

$$p_0 k - p_1 \frac{\gamma g Q A}{L} = 0, \quad p_1 k - p_0 \frac{\gamma Q A}{L} = 0.$$
 (5)

In the first equation in Eqs. (5), we omit the contribution from L-g non-neutral neighbors. In the second equation, we omit the contribution from the second Hamming class. Both these corrections are proportional to 1/g. Using the balance condition  $p_1g = (QA - 1)/(A - 1)$ , we solve Eq. (5) to obtain

$$k = \gamma Q A \frac{\sqrt{g}}{L}, \quad p_0 = \frac{QA - 1}{(A - 1)\sqrt{g}}, \quad p_1 = \frac{QA - 1}{(A - 1)g}.$$
 (6)

Thus,  $k \ll 1$ , which is consistent with Eq. (4).

The above calculations illustrate well that the impact of neutrality is determined by the maximal degree g, and its effect on the critical chain length is only of the order  $\sim \sqrt{g}/L < 1/\sqrt{L}$  because g < L. We can also derive the probability for a sequence with larger Hamming distances from  $S_0$ . For this purpose, let us assume now that a neutral mutant from the first class has a neutral neighbor from the second class with a relative probability  $p_2$ . It follows from Eq. (1) that  $p_2k = \gamma QAp_1/L$  or  $p_2 = p_1/\sqrt{g}$ , which gives

$$p_2 = \frac{QA - 1}{(A - 1)g^{3/2}}. (7)$$

Repeating the derivations for the probability of a neutral sequence at Hamming distance l (along the neutral network)  $p_l$ ,  $l \ge 1$ , we can show that the probability of having neutral sequences at the Hamming distance l along the neutral network from the master sequence is:

$$p_l = \frac{QA - 1}{(A - 1)g^{(1+l)/2}}. (8)$$

We can use results by Nimwegen et al [14] to check the reliability of our result. For a very small  $\gamma$ , Eqs. (4) and (6) in [14] by Nimwegen et al can be written as

$$p_i \bar{\mathbf{g}} = \sum_j G_{ij} p_j, \quad k \equiv R - Q A = \gamma Q A \frac{\bar{\mathbf{g}}}{L}, \tag{9}$$

where G is the adjacency matrix of the neutral network:  $G_{ij} = 1$  if types i and j are neighbors, and otherwise  $G_{ii} = 0$ . In the case of the Eigen model, considered in the current article, Eq. (9) is derived for the finite  $\gamma$  as well. Equation (9) could easily be solved for the types of neutral network shown in Figs. 3 and 4.

For the neutral network type of Fig. 3, we obtain

$$\bar{g} \approx \sqrt{g} + O(1/g),$$
 (10)

of which the derivation is presented as Case 5 in Materials and **Methods**. Equations (9) and (10) are consistent with Eq. (6).

For the neutral network type of Fig. 4 with a large degree g, we have

$$\bar{g} \approx \sqrt{g} + \frac{1}{2\sqrt{g}},$$
 (11)

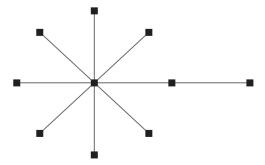
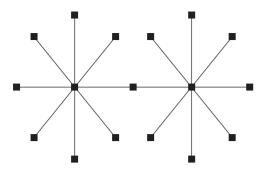


Figure 3. Neutral network-like fitness landscapes. There is a sequence with g=8 neutral neighbors and a tail of neutrals with the

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**Figure 4. Neutral network-like fitness landscapes.** There are two sequences with  $g\!=\!8$  neutral neighbors, with  $l\!=\!2$  distance between two centers.

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of which the derivation is presented as Case 6 in **Materials and Methods**. Thus our results are consistent with those obtained by Nimwegen *et al* [14].

As another test of Eq. (6), we present the derivation for the simplest case of one central sequence and its g neutral neighbors i,  $1 \le i \le g$ , as follows. Nonzero matrix elements of  $G_{ij}$  are  $G_{0i} = G_{i0} = 1$ . Due to the symmetry, Eq. (9) transforms into the system of equations:  $p_0\bar{g} = p_1g$ , and  $p_1\bar{g} = p_0$ . Multiplying the first equation by the second gives  $\bar{g} = \sqrt{g}$ . With such a result in the second expression of Eq. (9), we can obtain the first expression of Eq. (6). Thus our results are consistent with those obtained by Nimwegen  $et\ al\ [14]$ .

**Mesa-type fitness landscape.** The mean fitness of the *mesa* landscape has been calculated first in [21–23], then later analyzed in [28]. In a typical case, the high fitness extends to the Hamming distance of  $l_0$ , i.e.,  $J_{l \le l_0} = A$  and  $J_{l > l_0} = 1$ , where  $J_l \equiv r_i$  is the fitness at Hamming distance d = l from the wild sequence  $S_0$ , see e.g., Eq. (13) in [22] with  $l_0 = L(1-k_0)/2$  and  $k_0$  is defined in [22]. We derive rigorous solutions for several cases in "Correction terms for meta-type fitness landscapes" in **Materials and Methods**.

Solutions for various fitness landscapes are presented in Table 1. We note that because of the neutrality, there is a slight increase in the values of the mean fitness:  $QA \rightarrow QA(1+c\gamma/\sqrt{L})$ , where  $c \sim 1$ . Results for Case 3 and Case 4 in Table 1 show how by removing a single point, the  $1/\sqrt{L}$  correction term changes. This sensitivity of the mutant spectrum is quite typical. These results were derived by the  $1/\sqrt{L}$  expansion. Higher-order correction terms can be derived as well, i.e., the accuracy of the results included in Table 1 can be controlled.

Table 1 implies that the change in mean fitness and error threshold (defined by equation R=1) due to neutrality are rather small, of the order  $1/\sqrt{L}$ .

**Table 1.** Correction terms  $k \equiv R - QA$  (in the third row) for 4 different cases of mesa-type fitness, where R is the mean fitness and A is the fitness at the wild sequence  $S_0$ .

case	1	2	3	4
$r_l$	$r_0 = r_1 = A$	$r_1 = A$	$r_1 = r_2 = A$	$r_0 = r_1 = r_2 = A$
k	$\frac{\gamma QA}{\sqrt{L}}$	$O(\frac{1}{L})$	$\frac{\sqrt{2}\gamma QA}{\sqrt{L}}$	$\frac{\sqrt{3}\gamma QA}{\sqrt{L}}$

1.  $r_l \equiv A$  is the fitness at the Hamming distance d=l from  $S_0$ . 2.  $r_l=1$  for all  $r_l$  that are not explicitly included in the table. doi:10.1371/journal.pone.0021904.t001 The expression of Eq. (5) in [22] can be considered as the mean fitness R of the model defined by Eq. (1) in [22]. We have used Eq. (5) in [22] to discuss the effect of the flat fitness function, defined by Eq. (13) in [22]. It is easy to show that the modification of R is of the order  $A\gamma\sqrt{4l_0/L}$  for  $1-k\equiv 2l_0/L <<1$ . The  $1/\sqrt{L}$  result derived first in [22], was observed later in [28].

Neutral thick hierarchic tree-like fitness landscape with decreasing thickness of branches. The third scenario of neutrality is connected with a neutral thick hierarchic tree fitness landscape. It is a solvable model for the fitness landscape, where the genome is fractured into several parts with regard to the neutral property. For the Hamming distance l=1 from the peak configuration  $S_0$  with fitness A, there are (1-v)L sequences with fitness 1 and vL sequences with fitness a. At the Hamming distance a there are a there are a there are a to continued until the Hamming distance a from the peak configuration a to a the fitness a that Hamming distance larger than a always has the fitness a the Hamming distance larger than a always has the fitness a the Hamming distance exponentially until the maximal Hamming distance a.

We derive a close system of exact algebraic equations for any finite K, then check the quick convergence of the mean fitness shift with the K. We assume that such a model is close to the reality. In "Corrections for neutral landscape with thick hierarchic tree" in **Materials and Methods**, we calculated the corrections to the mean fitness. We find for the mean fitness the correction factor to be only of the order  $1/\sqrt{L}$ .

Following the results of [8,30] for one point neutral mutants we take the value v = 0.27; in "Corrections for neutral landscape with thick hierarchic tree" in **Materials and Methods** we further take K = 3 for our thick tree model to obtain

$$k = 0.549 \gamma Q A / \sqrt{L}. \tag{12}$$

The thick neutral-network fitness landscapes considered here change the error threshold only by a few percent, and certainly cannot solve the error-catastrophe paradox of the origin of life.

**Extensive neutrality.** Let us discuss different versions of extensive neutrality to clarify limits of applications for our non-extensive neutrality formulas. The main question we should address is, in which situations will the extensive neutrality change our result about the  $1/\sqrt{L}$  factor? First of all, we consider several situations with "thick" neutrality sub-manifolds (parts of neutral manifolds, connected with each other via neutral pathways), including partially mesa landscape, thick sub-manifold with a long thin tail, two thick sub-manifolds connected by a thin path, and two overlapping mesa landscapes.

In the section "Neutral selective Value" of [22], we have considered the "Partially mesa landscape". Consider some fraction  $\nu$  of alleles. Any mutations of these alleles with the total number d or less is neutral. We define such a landscape as a "partially mesa landscape".

For the increase of mean fitness due to neutrality in such landscapes, the left-hand side of Eq. [19] in [22] gives:

$$\delta R = \left[Ae^{\gamma(-1+\nu\sqrt{1-k_0^2})} - AQ\right] \sim 2\gamma QA\sqrt{\frac{\nu d}{L}}$$
 (13)

where we consider the case  $(1-k_0)/2 \equiv d/(vL) \ll 1$  and denote  $Q = \exp(-\gamma)$ . If the fitness does not depend on the values of nucleotides in some part of a genome with the length vL, then the error threshold changes  $\sim v$  because the effective genome length becomes shorter, L(1-v). In such a case of a partial mesa

landscape  $(k_0 = 0 \text{ in Eq. } (13))$ , there is no small factor  $(1/\sqrt{L})$ , the case considered in [6,16,17]). Equation (13) coincides with Eq.(4) in [17] for  $k_0 = 0$ .

It has been observed in experiments that there are rather long neutral pathways in a sequence space. Figure 3 is a schematic diagram of a landscape with long neutral paths in sequence space. But such long paths cannot significantly change the mean fitness, as has been well illustrated by Eq. (10) of the present paper. The long tail contribution to mean fitness is negligible compared with the "thick" part of the neutral manifold (supposed to be a partially mesa landscape), and the latter gives an increase of mean fitness  $O(\frac{1}{L})$ .

If we have several "thick" parts of a neutral manifold, connected together by thin neutral paths, the common increase of mean fitness due to neutrality is just equal to the increase by one with the maximal "thickness". Figure 4 is a schematic diagram for such a landscape with two "thick" parts of neutral manifold connected together with a thin neutral path. Equation (11) illustrates this phenomenon in the case of two identical "thick" neutral sub-manifolds, connected with the thin neutral path.

Let us consider overlapping mesa landscapes. In the simplest case, we have two reference sequences at the Hamming distance d, and the sequences are neutral until n mutations from either of the reference sequences.

We performed numerics for the parallel model [18], which is closely connected with the Eigen model [20]. In Table 2, we provide the results of the mean fitness for different distances between central sequences. We see only a slight increase in the neutrality impact O(1/L), and the maximal fitness increase appears at a small distance d=2. While we have done numerics only for two overlapping mesa landscapes, it is reasonable to assume that the same O(1/L) scale of corrections should still be valid in cases with several overlapping mesa landscapes.

The results listed in Table 2 are for the parallel model with  $\gamma = 1$ mutation rate per genome and J=2 difference between fitnesses of the sequences on the neutral manifold and the reminder sequences.

Experimental data analysis for the effect of extensive **neutrality.** Until now we have assumed a perfect neutrality, when the neutral mutants have exactly the same fitness as the master sequence  $S_0$ . Let us now analyze the experimental data of [8,30], to clarify the possible modification of our theoretical conclusions for imperfect neutrality, corresponding to observed data. The authors of [8,30] defined as "neutral" sequences having  $\epsilon = 0.038$  less relative fitness than the master sequence has. How large is such a decrease in fitness? Our formulas for the neutral network are valid when the nearest neighbors have a decrease in relative fitness  $\epsilon \ll \frac{\gamma}{\sqrt{L}}$ . Otherwise, when  $1 \gg \epsilon \gg \frac{\gamma}{\sqrt{L}}$ , Eq. (5) gives another result for the change in mean fitness, when we have a central sequence with fitness A and its g = vL neighbors with fitness  $A(1-\epsilon)$ :

**Table 2.** Numerically obtained mean fitness  $R_{num}$  for the parallel model with two overlapping mesa landscapes around two sequences, with Hamming distance d between central sequences, and with the maximal neutral mutation number n=5.

d	0	1	2	3	4	5
$R_{num}$	1.033287	1.033386	1.033395	1.033288	1.033287	1.033287

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$$k = \frac{\gamma^2 Q A v}{\epsilon L}.\tag{14}$$

Consider now the case of extensive neutrality: we assume that the multiple neutral mutations act independently, and thus the relative Wrightian fitness after n neutral mutations is [31]

$$A(1-\epsilon)^n. \tag{15}$$

Using Eq. (5) from [22], we again obtain Eq. (14). Thus the nearest neighboring neutral mutants make the bulk of the contribution to the increase in mean fitness due to neutrality, according to the data by [8,30].

All our formulas are for the selective phase where  $\gamma \sim 1$ . Putting L=10000,  $\epsilon=0.038$ ,  $\nu=0.27$ , we find that the error threshold is changed by only 0.1%. This result does not change even if we take into account the epistasis, observed in [30].

#### **Lethal Mutants**

The existence of lethal mutants is well established experimentally [8] and there have been several approximate results [32,33]. A rigorous investigation of the phenomenon started only recently. In [9], we calculated the exact mean fitness for the model with a general symmetric fitness landscape and lethal mutations, including the case of the single peak landscape as a special case. The exact error threshold for the latter case was derived by Tejero et al. [35], who also used approximate methods and ideas of [36] to study the extinction threshold.

The extinction phenomena in bacteria originate from "internal" degradation: a mother bacterium is replaced by two daughter bacteria (with possible mutations), therefore the mother disappears after the self-replication cycle. In contrast to the case of bacteria, we assume that the self-replicating RNA molecule does not disappear after providing copies, and therefore can participate in self-replicating events multiple times. Therefore, there is neither an "internal" degradation process nor connected extinction threshold phenomenon in our case; see the subsection Extinction threshold in Materials and Methods. Here we will calculate the probability distribution for a single peak fitness model with lethal mutations, which was not done in [34] and [35].

Let us consider a single-peak fitness model, in which lethal mutants are randomly distributed in the sequence space. First of all we define accurately the distribution of lethal mutants (zero fitness) in the sequence space. The number M of non-lethal sequences scales as some degree of the total sequence number,

$$M \sim (2^L)^c, \quad 0 < c < 1.$$
 (16)

How can we dilute the sequence space by lethal sequences? Let us choose a reference sequence  $S_0$  (the sequence with a high fitness Ain the case of a single peak fitness model). In the first Hamming class with Hamming distance d = 1 from  $S_0$ , we have  $L_1 \equiv (1 - \lambda)L$  nonlethal mutations and  $\lambda L$  lethal mutations; in the second class with d=2, there are  $L_2 \equiv L(1-\lambda)[L(1-\lambda)-1]/2! = L_1(L_1-1)/2!$ non-lethal sequences;...; in the n-th class with d=n, there are  $L_n \equiv L_1!/[(L_1-n)!n!]$  non-lethal sequences. Thus the total number of non-lethal sequences is

$$M = \sum_{n=0}^{L_1} \frac{L_1!}{[(L_1 - n)!n!]} = 2^{L(1 - \lambda)}.$$

Comparing this M with that of Eq. (16), we have  $c = (1 - \lambda)$ . In the above derivation, we assume that any sequence having several mutations, including at least one lethal mutation from the onepoint lethal mutation list is also lethal, and we ignore the combinations of deleterious mutations (synthetic lethal). Such a picture is quite realistic for RNA viruses [8,30].

For the lethal mutants with parameter  $\lambda$ , in the subsection Lethal mutants in Materials and Methods, we derive with 1/L accuracy

$$p_0 = Q^{\lambda} \frac{Q^{1-\lambda} A - 1}{A - 1}.$$
 (17)

In the infinite population limit, the error threshold can be determined by the condition

$$p_0 > 0$$
.

Thus from  $p_0$  in Eq. (3), one can obtain the error threshold in Eq. (2) derived by Eigen. In the infinite population limit,  $p_0$  in Eq. (17) implies that the error threshold for the case with fraction  $\lambda$  of lethal mutants is given by

$$AQ^{1-\lambda} = 1. \tag{18}$$

Current experiments suggest that the probability  $\lambda$  of one point lethal mutants is  $\lambda \sim 40\%$ , i.e., about twice as high as the probability of hitting the neutral mutants [6,8]. For  $\lambda = 40\%$ , the error threshold constraint is relaxed by a factor of 5/3. This is insufficient to solve the error threshold paradox. It could be solved by increasing the degree of lethality  $\lambda$  and involving the truncated selection (see next section). Equation (18) shows that while the lethal mutations change the error threshold, the fraction of the master sequence decreases with the high mutation rates (small Q). If there is an extinction threshold in the population (the population disappears below a minimal value of the mean fitness) [35–37], then even the lethal mutations cannot rescue the situation: the selective phase disappears.

The existence of the error threshold is a fundamental phenomenon, connected with the Shannon optimal codes in information theory [38], while the extinction threshold is a case dependent, non-universal phenomenon.

#### Truncated selection with lethal mutations

Consider a fitness landscape with "truncated selection" (Fig. 2) [10-12]. We take  $r_0 = A$ ;  $r_l = 1$ ,  $1 \le l \le n$ ; and,  $r_l = 0$  for l > n, where n is the truncation parameter. We denote by  $p_l$  the probability of having a sequence from the l-th class. In [12], we solve analytically the truncated fitness landscape for the case of large n, and perform numerics for the finite n case. For the large n, we find in [12] that the error threshold transition is fractured into two separate transitions.

Now we will derive analytical expressions of the mean fitness, and consider the case of truncated mutation with lethal mutants in the case of small n. For the master type, we have  $p_0(QA - R) = 0$ ,

where  $R = Ap_0 + \sum_{l=1}^{n} p_l \frac{L^l}{l!} = p_0 (A + \sum_{l=1}^{n} z_l \frac{\gamma^l}{l!})$ , where  $z_l$  are defined by Eq. (54) in **Materials and Methods**. Then we define the  $p_0$  by the equation

$$p_{0} = \frac{QA}{A + \sum_{l=1}^{n} z_{l} \frac{\gamma^{l}}{l!}} = \frac{QA}{\phi_{n}(\gamma, A)},$$
(19)

where the function  $\phi_n$  is defined in Eq. (55) in **Materials and Methods**. For the truncated selection in the presence of lethal mutants, Eq. (19) should be changed to:

$$p_0 = \frac{QA}{A + \sum_{l=1}^{n} z_l \frac{\gamma^l (1 - \lambda)^l}{l!}} = \frac{QA}{\phi_n(\gamma(1 - \lambda), A)}.$$
 (20)

We put the error threshold condition within the 1/L accuracy:

$$p_0 \ge 1/L. \tag{21}$$

Were a population size N available from experiments, we could, instead of Eq. (21), use another constraint for the  $p_0$ :

$$p_0 N \gg 1. \tag{22}$$

The justification for the conditions given in Eq. (22) is that molecular population size N should be high enough to produce deterministic features. The population size should be large enough to avoid the loss of the master sequence due to the Muller's ratchet effect. In reality, it is not easy to obtain the value of N. Thus in the following, we will use Eq. (21) to estimate the critical length.

Having  $p_0$  given by Eqs. (17) or (20) for several typical parameters of the models, our estimates of the critical length with Eq. (21) are gathered in Table 3. The results for the neutral network case and neutral thick network case were obtained with the use of Eq. (6) with g = vL, and Eq. (12), respectively. The degree of neutral mutations is taken as  $v \approx 0.27$  [8]. We have verified that for the four-letter alphabet the impact of neutrality is smaller by a factor of  $\sqrt{3}$ , but other entries in Table 3 remain unchanged.

If we assume  $v = 0.1, \lambda = 0.40$ , and extensive neutrality according to Eq. (15), then the results of Table 3 are changed slightly by 1.5 - 2%.

#### Discussion

In this work we have rigorously investigated the error-threshold problem for evolution with neutral and lethal mutants, and with truncated selection. We have calculated the change in mean fitness (e.g. k in Eqs. (4) and (6)) due to neutrality for the neutral network with a high degree g at some node. We also considered the neutrality phenomenon for a more involved case, in which the fraction of neutral mutants among all multiple mutants decreases exponentially with the Hamming distance from the master sequence. Then we found that the neutrality changes the mean fitness and the error threshold by only a few percent  $1/\sqrt{L}$  and certainly cannot solve Eigen's error threshold paradox. The formulas considered correspond to perfect neutrality, where neutral mutants have exactly the same fitness as the master sequence. We also considered the case of imperfect neutrality: the neutrality case according to the data of [8], in which there is a small decrease in fitness after mutations. Assuming extensive

**Table 3.** The maximal allowed genome length  $L_{max}$  obtained from different conditions (equations) for several values of the parameter of truncated selection n and the degree of lethal mutations  $\lambda$  when  $\ln A = 1$  and 1 - q = 0.001.

Conditions	$L_{max}$	n	λ
Eq. (2)	1000		
Eqs. (3), (21)	998		
Eq. (6) with 27%-neutrality	1016		
Eq. (12) with 27%-neutrality	1017		
Eq. (18)	1666		0.4
Eqs. (17),(21)	1664		0.4
Eqs. (17),(21)	2000		0.5
Eqs. (17),(21)	4915		0.8
Eq. (18)	5000		0.8
Eqs. (17),(21)	6300		0.85
Eq. (18)	6666		0.85
Eqs. (17),(21)	7800		0.9
Eq. (18)	10000		0.9
Eqs. (20),(21)	4650	4	0.4
	5430	3	0.4
	6500	2	0.4
	5050	4	0.5
	5800	3	0.5
	6750	2	0.5
	7050	4	0.8
	7450	3	0.8
	7900	2	0.8
	8200	3	0.9
	8310	2	0.9

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neutrality in such a case (multiplicative character of fitness for those mutations), we find that such a neutrality alone can change the error threshold by  $\sim 0.1\%$ . Thus neutral mutation alone cannot solve the error threshold paradox. Our result is rather general, as, according to experimental data [8], an assumed extensive neutrality gives the same increase in mean fitness as the model with only the nearest neighboring neutrals.

We solved exactly the model with lethal mutations. Both phenomena, the occurrence of significant proportions of the neutral and the lethal types, suppress the error threshold in a similar way, while there is neither a small factor  $1/\sqrt{L}$ , nor a finetuning problem in the case of lethal mutations. The effect of the lethal mutants is, however, easier to realize than the effect of neutrality, even after a billion-year evolution. The difference in the impacts of these two mechanisms is thought to have been even more pronounced at the origin of life. We have provided evidence that in modern RNA, the presence of lethal mutants can cause an increase in the error threshold by as much as 50%.

We have showed that the Eigen's error threshold for the origin of life can be relaxed, provided the presence of the lethal mutants is aided by truncated selection (see Table 3). For example, in the case of RNA molecules, the maximum length of chains is considerably extended when lethal mutants with an 80% lethal probability are included in the model, together with truncated selection. In the absence of truncated selection the probability of the master-type sequence would be negligible, which in turn would require enormously large molecular populations for reactions to happen. For maintaining a continuous replication, it is important to have in the population both lethal mutants and viable mutants. Moreover, the latter should be restricted to 2 to 4 base exchanges. If we assume 80% lethal probability and 10% neutrality, then the neutrality can change the error threshold by 1.5% - 2%.

One of the questions of interest concerns the organization of truncated selection in pre-biotic evolution. Only recently it has been realized that proteins are not random heteropolymers but their sequences are formed following a tentative design (for a review see [39,40]). Developed during evolution, this design entails, for example, the robustness of the genome against mutations. In the context of applications to RNA, the concept of design was recently studied by Zorn et al [41]. It is reasonable to assume that the degree of design and the robustness were poor at the beginning of evolution. For example, initial evolution might have followed a scenario in which the truncated selection took place in a population with a large number of lethal sequences. As shown here, in this example under poor organization, the error catastrophe could have been avoided.

The key point of our study is that no matter where the beginning of life was (it is an obligatory property of the matter, as has been assumed in [42]), if it was through autocatalytic reactions, it had to be accompanied by lethal mutants with truncated selection.

In summary, in populations that contain about  $\sim 80\%$  of lethal mutants and provide for the simultaneous truncated selection with the truncation parameter n=2 or 3, the primordial genome can reach the critical length of 7000 ~ 8000 estimated by Kun, et al. [6], and Eigen's error catastrophe can thus be solved.

There are three essential stages in the origin of life [42]. The First is the preliminary stage, with the preparation of the proper bio-molecules for the starting point [42-44]. The Second is connected with self-replication of macro-molecules [42]. The Third gives the protocells [45]. The present paper studies the second stage. The error threshold problem exists for both the second and third stages. In the second stage there is a replication of molecules using a template, while in the third stage the mother protocell divides into two protocells. The mechanism we suggested solves the error threshold for the second stage, but not for the third stage: too much lethal mutation push the population to the extinction threshold [15] and the self-replication of proto-cell will stop. Thus a protocell should have auto-proof mechanism of selfreplication to suppress the mutation rates.

Let us briefly discuss our results in view of alternative ideas to solve the error threshold paradox. All of our derivations and conclusions correspond to the case of replication of a pre-biotic molecule using a template. An alternative mechanism to avoid the error threshold could be connected with the self-replication of the network of molecules [46,47], in which several enzymes catalyze the generation of each other. Such a mechanism increases the value of the joint fitness of the "peak" configuration of a couple of sequences, which is useful in avoiding the error catastrophe. For the origin of life, we need some minimal pool of genes, which could be provided by two molecules (replicating together) with a shorter length for each chain. Unfortunately, the information contents of the two sequences in [47] are almost identical, therefore such a concrete mechanism could not provide a larger number of genes than the single sequence. In the case of a connected replication of several RNA-like molecules with different information content and self-sustaining amplification of the whole molecular group, such a mechanism, combined with the lethal mutations, can easily solve the error paradox, and the mathematical tools developed in the current article could be applied in this case as well.

Peck and Waxman [48] proposed the evolution model with recombination and concluded that the truncated selection and recombination could solve the error paradox. We agree with the importance of the truncated selection, while have not see serious argumentation for the importance of recombination to solve the error threshold catastrophe. They used non-zero degradation in their model, while have forgotten to analyze the extinction threshold. A rigorous consideration of the single peak fitness landscape with the simplest version of recombination in [49] proves that the recombination does not change the (mean fitness) error threshold for the long genome and hence could not solve the error paradox. For the short genome the recombination even slightly suppresses the selection (the mean fitness decreases) for the single peak fitness case [50].

Rajamani, et al. [51] have considered the mechanism of selfreplication cycle in details, assuming slow reaction rates for the mutants due to "mismatch stalling", which can somehow change the error threshold, when the error probability per nucleotide times the "stalling" coefficient [51] is larger than the fitness ratio (wild sequence fitness to the other non-lethal sequence fitness). Actually the considered phenomenon is equivalent to some increasing of the fitness ratio A. The phenomenon depends on the concrete details of the self-replication cycle. One should consider this phenomenon together with lethal mutations, truncated selection and finite period of generation [52].

Our work helps to solve a puzzle in the second stage of the origin of life [42]. Such result and other recent advances in models of cells [53-56] and minivirus [57] will provide clues for understanding the evolution from the second stage to the third stage of the origin of life.

## **Materials and Methods**

#### Correction terms for mesa-type fitness landscapes

Consider the steady state solutions of the Eigen model [1,2] for the fitness landscapes with two classes of sequences: with a high fitness A and with a lower fitness 1 (one).

**Single peak fitness model.** Consider the fitness landscape

$$r_0 = A, \quad r_l = 1, l > 0.$$
 (23)

This gives the following expression for the  $p_0$  and mean fitness [22]

$$p_0 = \frac{QA - 1}{A - 1},$$

$$R = QA. (24)$$

**General case.** Consider the fitness landscape where there is a high fitness A for  $l_1 \le l \le l_2$  and fitness 1 for other sequences. There is an exact equation for the mean fitness

$$R = 1 + (A - 1) \sum_{l_1 \le l \le l_2} p_l N_l, \tag{25}$$

where  $N_l = \frac{L!}{l!(L-l)!} \sim \frac{L^l}{l!}$  for small  $l \ll L$ . The total probability of neutral sequences is approximately the same as in the single peak fitness model,

$$\sum_{l_1 \le l \le l_2} p_l N_l - \frac{QA - 1}{A - 1} = O(\frac{1}{\sqrt{L}}). \tag{26}$$

We assume the following ansatz for  $p_l, l_1 \le l \le l_2$ 

$$p_l \sim \frac{1}{L^{l_2 - (l_2 - l)/2}}$$
 (27)

The mean fitness R is almost the same as for SP case:

$$R - QA \sim O(\frac{1}{\sqrt{L}}). \tag{28}$$

The majority of population is in the highest Hamming class and we have

$$p_{l_2} = \frac{(QA - 1)}{A - 1} \frac{l_2!}{L^{l_2}}.$$
 (29)

Then from Eq. (1), we have the following system of equations

$$kp_l = l\frac{\gamma QAp_{l-1}}{L} + \gamma QAp_{l+1} \tag{30}$$

where k = R - QA. For  $l = l_1$  we miss the first term  $\sim p_{l_1-1}$  and for the  $l = l_2$  we missed the second term  $\sim p_{l_2+1}$  in Eq. (30). The higher terms  $\sim p_{l+2}$  as well as the lower one  $\sim p_{l-2}$  are missed, because their contribution are suppressed due to a small factor  $1/\sqrt{L}$ .

Case 1. The simple mesa. Consider now the case, when besides the 0-th configuration there is a high fitness at the Hamming distance 1

$$r_0 = A \cdot r_1 = A$$
.

$$r_l = 1, l > 1.$$
 (31)

From Eq. (29), we have  $p_1L = \frac{QA-1}{4-1}$ . Equation (30) implies that

$$kp_0 = \gamma Q A p_1, kp_1 = \frac{\gamma Q A p_0}{I}, \tag{32}$$

where k = R - QA. We have a solution:

$$k = \frac{\gamma Q A}{\sqrt{L}}, p_0 = \frac{Q A - 1}{A - 1} \frac{1}{\sqrt{L}}.$$
 (33)

Case 2. Mesa with the hole. Consider now the case, when there is high fitness at the first Hamming class

$$r_0 = 1, r_1 = A, J_l = 1, l > 1,$$
 (34)

We have in the bulk approximation

$$p_1 = \frac{\gamma(QA - 1)}{L(A - 1)}, p_0 = \frac{\gamma A}{(A - 1)}, p_1, p_2 = \frac{2\gamma A}{(A - 1)L}, p_1.$$
 (35)

As  $Lp_1 \sim 1$ , then  $p_0 \sim 1/L$ , therefore there are no  $\frac{1}{\sqrt{L}}$  corrections now, just O(1/L) ones. To get k we consider the equation for  $p_1$ with the small corrections:

$$(R - QA)p_1 = \frac{\gamma Q}{L}p_0 + \gamma Qp_2, \tag{36}$$

which gives

$$k = R - QA = \frac{3\gamma^2}{L(A - 1)}. (37)$$

Case 3. Band with zero at the center. Consider now the fitness landscape

$$r_0 = 1, r_1 = A, r_2 = A, r_l = 1, l > 2.$$
 (38)

Now we have

$$p_1 k = \gamma Q A p_2, p_2 k = \frac{2\gamma Q A}{L} p_1.$$
 (39)

We derive immediately

$$k = \frac{\sqrt{2}\gamma QA}{\sqrt{L}}. (40)$$

Case 4. Thick band. Consider now the case, when

$$r_0 = A, r_1 = A, r_2 = A, r_l = 0, l > 2.$$
 (41)

We have equations

$$p_0 k = \gamma Q A p_1$$

$$p_1k = \frac{\gamma QAp_0}{L} + \gamma QAp_2,$$

$$p_2 k = \frac{2\gamma Q A p_1}{I},\tag{42}$$

which imply

$$k = \frac{\sqrt{3}\gamma QA}{\sqrt{L}}. (43)$$

The results of Cases 1-4 are listed in Table I.

Case 5. Model with fitness by Fig. 3. Denote the number of all one point neutral mutants by g,  $x = p_0$ , and by y the probabilities of g-1 similar neutral neighbors (without neutral tails), by z the probability for non-symmetrically located one point mutation neighbor of the master sequence, and by h the probability of two point mutation neutral mutant. We have a system of equations for variables x, y, z, h and the average number of degree for the whole neutral network  $\bar{g}$ .

$$\bar{g}x = (g-1)y + z, \bar{g}y = x, \bar{g}z = x + h, \bar{g}h = z.$$
 (44)

Putting  $h=z/\bar{g}$ ,  $y=x/\bar{g}$ , we derive an equation for  $\bar{g}$ :

$$\bar{\mathbf{g}}^2(\bar{\mathbf{g}}^2-1)-(g-1)(\bar{\mathbf{g}}^2-1)-\bar{\mathbf{g}}^2=0.$$
 (45)

For the large g we have

$$\bar{g} \approx \sqrt{g} + O(\frac{1}{g})$$
 (46)

If we take the longer tail, l>2 in the Fig 3, only O(1/g) terms are affected in Eq.(46).

Case 6. Model with fitness by Fig. 4. Denote again by g the total number of neutral neighbors of two nodes with the largest degree,  $x = p_0$ , and by v the probabilities of g - 1 similar neutral neighbors, by z the probability for non-symmetrically located one point mutation neighbor of the master sequence. We have a system of equations for x, y, z, and  $\bar{g}$ 

$$\bar{g}x = (g-1)y + z, \bar{g}y = x, \bar{g}z = 2x$$
 (47)

Putting  $z = 2x/\bar{g}$ ,  $y = x/\bar{g}$ , we derive an equation for  $\bar{g}$ :

$$\bar{\mathbf{g}}^2 = g + 1 \tag{48}$$

For the large d we have

$$\bar{g} \approx \sqrt{g} + \frac{1}{2\sqrt{g}} \tag{49}$$

We took the Hamming distance 2 between two centers of thick sub-manifolds. If we take more Hamming distance, then the second term in the last equation should be O(1/g).

## Corrections for neutral landscape with thick hierarchic tree

Here we follow the terminology of Sec. I to call a nucleotide as a spin. In the model of neutral landscape with thick hierarchic tree, the L spins are partitioned into K+1 different groups, with  $Ll_n$ spins in the *n*-th group, where  $l_n = (1-h)h^{n-1}v, 1 \le n < K$ , and  $l_K = vh^{K-1}$ . Thus the maximal distance on the tree from the reference sequence equals K. There are at most n neutral pointmutations in the *n*-th group. We take h=v. Compared with the neutral network of the previous subsection, now the hierarchic tree has a thick tube instead of thin lines. What we are doing is equivalent to solving Eq. (9) using the symmetry: collecting the same type of sequences together. We should distinguish among different probabilities for the sequences that are obtained from the central sequence after mutations of the spins from different groups. We denote as  $X_{ln}$  the probability of having a sequence from the *l*-th Hamming class but with the spins from the *n*-th group. Such sequence has l''-l'' spins and  $l_nL-l''+l''$  spins from the same group of spins. From the Eigen model equations we get:

$$X_{l,n}k = \frac{QA\gamma}{L}(lX_{l-1,n} + (Ll_n - l)X_{l+1,n}),$$

$$0 \le l \le n, 1 \le n \le K. \tag{50}$$

and we identify  $X_{0,n} = X_0$  and  $k \equiv R - QA$ , also  $X_{-1,n} = 0, X_{n+1,n} = 0.$ 

We have a complete system of equations to define k and  $X_{ln}/X_0$ . In the case when  $Ll_K > > 1$ , we can further simplify the system of equations,

$$X_{l,n}k = \frac{QA\gamma}{L}(lX_{l-1,n} + Ll_nX_{l+1,n}).$$
 (51)

Consider first the case of K=3. We have  $X_0$  for the 0-th class probability. For the lowest group we have  $X_{11}$  for a sequence in the first Hamming class, i.e., when the distance from the master type is d=1. For the second group of sequences, i.e., for those obtained via mutations of the second spin group, we have  $X_{12}$  in the first Hamming class, and  $X_{22}$  in the second Hamming class. For the third group of sequences, we have  $X_{13}$  in the first Hamming class, and  $X_{23}$  in the second Hamming class, and  $X_{33}$  in the third class. We derive the following system of equations:

$$k_0 X_{11} = X_0 / L, k_0 X_0 = l_1 X_{11} + l_2 X_{12} + l_3 X_{13},$$

$$k_0 X_{12} = X_0/L + l_2 X_{22}, k_0 X_{22} = 2X_{12}/L,$$

$$k_0 X_{13} = X_0/L + l_3 X_{23}, k_0 X_{23} = 2X_{13}/L + l_3 X_{33},$$

$$k_0 X_{33} = 3X_{23}/L. (52)$$

Here  $k_0 = k/(\gamma QA)$  is an eigenvalue of a matrix. The second equation in the first line of Eq. (52) was derived directly from Eq.

After re-scaling  $k_0 = \lambda/\sqrt{L}$ ,  $X_{ln}/X_0 = y_{ln}/L^{l/2}$  we have a system of equations without large parameter L:

$$\lambda y_{11} = 1, \lambda = l_1 y_{11} + l_2 y_{12} + l_3 y_{13},$$

$$\lambda v_{12} = 1 + l_2 v_{22}, \lambda v_{22} = 2v_{12},$$

$$\lambda y_{13} = 1 + l_3 y_{23}, \lambda y_{23} = 2y_{13} + l_3 y_{33},$$

$$\lambda y_{33} = 3y_{23}. (53)$$

In the last equation  $l_n$  are functions of v. For v = 0.27 we have a correction to the mean fitness  $0.549\gamma QA/\sqrt{L}$ . The K=3 gives the results of large K with the accuracy 0.1%.

## Lethal mutants

We will calculate mean fitness with 1/L accuracy. Let us first consider the case without lethal mutants. We denote by  $p_l$  the probability of having a sequence from the l-th class, and by  $P_l$  the probability of having the l-th class. In the l-th Hamming class there are  $N_l = \frac{L!}{l!(L-l)!} \approx L^l/l!$  sequences and  $P_l \equiv N_l p_l$ . Let us

denote  $L^l p_l / (p_0 \gamma^l) \equiv z_l (\gamma, A)$ , and therefore  $z_0 = 1$ . The recurrence relations for  $z_l$  are [22]:

$$z_{l} = \frac{1}{A - 1} \left[ A + \sum_{1 \le k \le l} \frac{z_{k} l!}{k! (l - k)!} \right].$$
 (54)

Having the values of  $z_l$  we can calculate  $p_l = p_0 z_l \gamma^l / L^l$ .

We have for the mean fitness R = QA. On the other hand, mean fitness is defined as  $R = Ap_0 + \sum_{l=1}^{L} p_l \frac{L^l}{l!} + O(1/L)$ . Thus with the O(1/L) accuracy,

$$p_0 = \frac{QA}{\left[A + \sum_{l=1}^L \frac{z_l \gamma^l}{l!}\right]} \equiv \frac{QA}{\phi_L(\gamma, A)}.$$

Using the master degree probability  $p_0 = \frac{QA-1}{A-1}$ , we get,

$$\phi_L(\gamma, A) \equiv A + \sum_{1=l}^{L} \frac{z_l \gamma^l}{l!} = e^{-\gamma} A \frac{A-1}{e^{-\gamma} A - 1}.$$
 (55)

Consider now the case with lethal mutants. In the l-th Hamming class we have  $N_l(1-\lambda)^l$  non-lethal sequences: each of which has the probability  $x_l$ , and  $N_l[1-(1-\lambda)^l]$  lethal sequences. For the  $x_1/x_0, x_0 = p_0$  we have the same system of recurrent equations as those for  $p_I/p_0$  in case without lethals, therefore we can calculate  $x_l$  as  $p_l$ . A single modification, we should replace  $N_l \rightarrow N_l (1-\lambda)^l$  while calculating the mean fitness expression:

$$p_0 A + \sum_{l=1}^{L} N_l (1-\lambda)^l x_l \approx p_0 \phi(\gamma(1-\lambda), A)$$

$$=p_{0}e^{-\gamma(1-\lambda)}A\frac{A-1}{e^{-\gamma(1-\lambda)}A-1}.$$
 (56)

For the distribution of lethal sequences with the faction  $\lambda_l$  in the *l*-th class [33], the last equation is modified:  $(1-\lambda)^l \rightarrow \lambda_l$ .

Equation (56) defines the mean fitness. For the master type we have  $p_0(QA-R) = O(1/L)$ . Thus we have again for the mean fitness  $R = QA = e^{-\gamma}A$ . Therefore, we have an equation for the  $p_0$ 

$$e^{-\gamma}A = p_0 e^{-\gamma(1-\lambda)} A \frac{A-1}{e^{-\gamma(1-\lambda)}A-1}$$

which implies Eq. (17).

#### Lethal and neutral mutants

Consider a single peak sequence with a fitness A, a part of genome  $\nu L$  with a fitness by Eq. (15), and  $\lambda L$  positions in genome for lethal mutations. In the selective phase we get, following to Eq.(14), a mean fitness

$$(1 + \frac{\gamma^2 Q A v}{\epsilon L}) Q A. \tag{57}$$

Using Eq. (18) and the expression QA for the mean fitness of selective phase, we get the following expression for the mean fitness of non-selective phase:



$$Q^{\lambda}$$
. (58)

Comparing the latter two expressions, we get for the error threshold:

$$(1 + \frac{\gamma^2 Q A \nu}{\epsilon L}) A Q^{(1-\lambda)} = 1. \tag{59}$$

## Extinction threshold

The growth of bacterial population is through the cell division. At time t, we have  $n_j(t)$  bacteria of the given type. Let us consider the division of the bacteria of type i into two daughter bacteria with types k and l [37]. After the bacteria division the number of bacteria of the type i decreases for 1, and with probabilities  $Q_{ik}, Q_{il}$  increases the numbers of bacteria with the types k, l.

$$n_i(t+1) = n_i(t) - 1$$
,

$$n_k(t+1) = n_k(t) + 1$$
,

$$n_l(t+1) = n_l(t) + 1.$$
 (60)

We can model such a situation with the continuous time model

$$\frac{dp_i}{dt} = \sum_k Q_{ik} r_k p_k(t) - d_i p_i(t), \tag{61}$$

where  $d_i$  is the degradation term introduced in [2], and following to [36] and [37]. All the terms  $Q_{ij} \sim e^{-\gamma}$ . The origin of the  $d_i$  is just the disappearance of the mother bacteria after the division ("internal" degradation).

Carefully analyzing Eq. (61), the authors of [36,37] deduced the extinction threshold, a phenomenon when the total population size

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decreases. There is a strict constraint

$$2e^{-\gamma} > 1,\tag{62}$$

otherwise the population disappears.

The origin of their conclusion is the existence of nonzero degradation rate  $d_i$ , initiated by the first equation in Eq.(60).

The point is that in the case of RNA replication considered in this paper, we have another situation:

$$n_i(t+1) = n_i(t)$$

$$n_k(t+1) = n_k(t) + 1.$$
 (63)

Thus while considering the corresponding continuous time model, we don't need to add the negative term, and get just the Eigen model for growing population

$$\frac{dp_i}{dt} = Q_{ii}r_ip_i(t) + \sum_{k \neq i} Q_{ik}r_kp_k(t)$$
(64)

Of course, it is possible some degradation due to interaction with external environment, but there is no strict constraint like the one described in Eq. (62).

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# **Author Contributions**

Conceived and designed the experiments: DBS CKB CKH. Performed the experiments: DBS CKH. Analyzed the data: DBS CKH. Contributed reagents/materials/analysis tools: DBS CKH. Wrote the paper: DBS CKH.

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